

MPDATA i libmpdata++:

O jawnym schemacie numerycznym wysokiego rzędu
całkowania układów wielowymiarowych równań transportu
i jego nowej implementacji w C++

Sylwester Arabas

- MPDATA (Smolarkiewicz '83 ... Smolarkiewicz et al. 20XX)
- libmpdata++ i przykłady z geofizyki (Jaruga et al. 2015)
- zastosowanie w finansach (Arabas & Farhat, arXiv)

MPDATA

transport PDE:
$$\frac{\partial \psi}{\partial t} + \frac{\partial}{\partial x} (v\psi) = 0$$

MPDATA in a nutshell (Smolarkiewicz 1983 MWR ...)

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$$F(\psi_L, \psi_R, \mathcal{C}) = \max(\mathcal{C}, 0) \cdot \psi_L + \min(\mathcal{C}, 0) \cdot \psi_R$$

$$\mathcal{C} = v\Delta t / \Delta x$$

← upwind

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$$C'_{i+1/2} = (|C_{i+1/2}| - C_{i+1/2}^2) A_{i+1/2}$$

$$A_{i+1/2} = \frac{\psi_{i+1} - \psi_i}{\psi_{i+1} + \psi_i}$$

MPDATA: reverse numerical diffusion by integrating the antidiffusive flux using upwind (in a corrective iteration)

MPDATA: key features (review: e.g. Smolarkiewicz 2006)

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- ❖ **ERC-winning:**
Piotr’s PantaRhei Advanced Grant @ ECMWF

Modeling the Solar Dynamo

Paul Charbonneau¹ and Piotr K. Smolarkiewicz²

The Sun's magnetic field is the engine and energy channel underlying virtually all manifestations of solar activity. Its evolution takes place on a wide range of spatial and temporal scales, including a prominent 11-year cycle of successive polarity reversals over the entire star. This magnetic cycle in turn modulates the physical properties of the plasma flowing away from the Sun into interplanetary space, the frequency of all geoeffective eruptive phenomena (such as flares and coronal mass ejections), and the solar radiative flux over the full range of the electromagnetic spec-

trum—from x-rays through ultraviolet, visible, and infrared light, all the way down to radio frequencies (1). The Sun's heartbeat is truly magnetic, and recent numerical simulations (2–5) are providing new insights into its mode of operation.

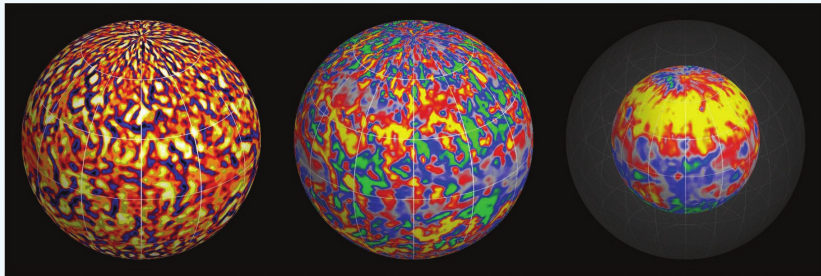
Self-sustained amplification of a magnetic field through the action of fluid motions is called a dynamo. Dynamos operate through a physical effect called electromagnetic induction, discovered by Faraday in the 19th century. Induction is put to work in modern power-generating plants converting mechanical energy imparted to turbines (by water, wind, or steam) into electricity. There are no turbines inside the Sun, but in its outer third in radial extent, the so-called convection zone, mechanical energy abounds in the form of rotational shear and

Numerical simulations are changing our views on the dynamo process underlying the solar magnetic activity cycle.

turbulent fluid motions driven by the solar luminosity. Plasma flowing across the magnetic field that pervades the solar interior induces electrical currents, which, under appropriate flow and magnetic field configurations, can sustain the field against dissipation. The magnetic field so generated in the solar interior subsequently emerges at the Sun's surface, structuring and energizing its extended atmosphere.

Parallel advances in raw computing power and ever more sophisticated numerical algorithms make it possible to produce and investigate magnetic cycles in Sun-like spheres of thermally convecting magnetized fluid. The underlying physics is in principle well understood in the form of magnetohydrodynamics (or MHD), being described by the classical fluid equations augmented by

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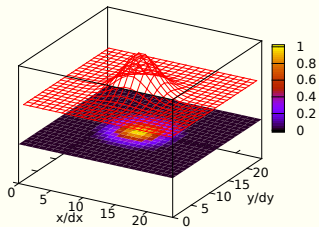


Solar simulations. Three snapshots of a magnetohydrodynamical numerical simulation of solar convection, carried out using the multiscale flow simulation model EULAG (12–15). The left panel shows a color rendering of the radial component of the convective flow (orange to light yellow, upflows; red to dark blue, downflows) in the subsurface layers of the simulation. The center panel shows the radial magnetic field (gray to yellow, outward-directed magnetic field; gray to green, inward-directed) at the same depth. Note how the characteristic spa-

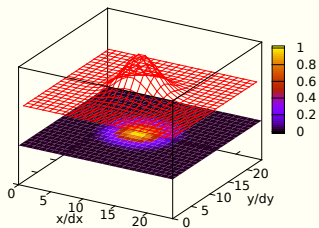
tial scales are the same for both quantities, which also evolve on the same temporal scale of days. The right panel shows the magnitude of the zonal magnetic field component, deep in the interior of the simulation, at the base of the convection layer. Note the banded structure at mid-latitudes, roughly symmetric about the rotation axis. This torus-like structure, and its opposite-polarity counterpart in the other hemisphere, undergo synchronous polarity reversals on a cadence of about 40 years.

przykład 2D (Arabas et al. 2014, Sci. Prog.)

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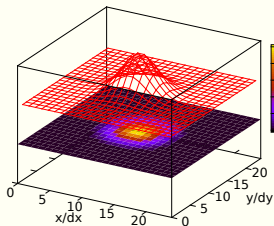


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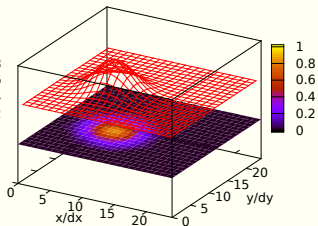


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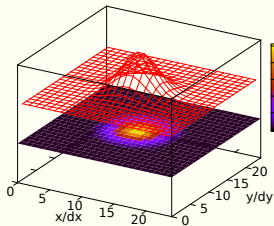
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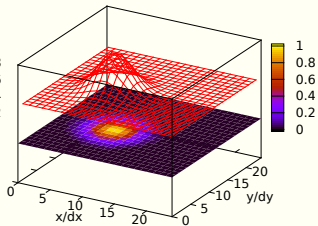
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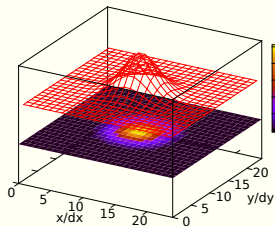


mpdata<3> $t/dt=6$

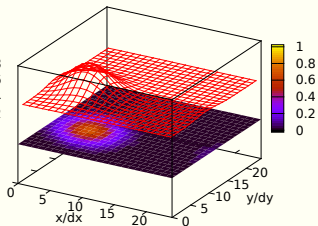


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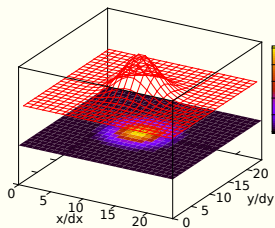
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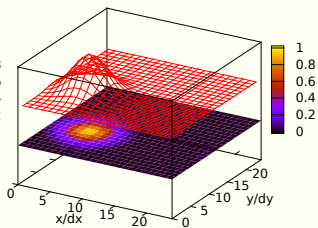
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mpdata<3> $t/dt=0$

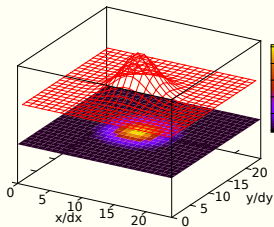


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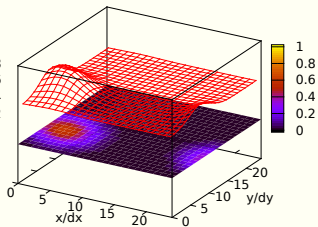


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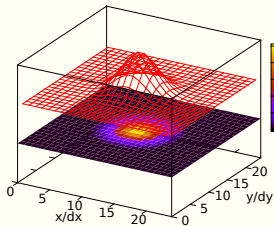
donorcell $t/dt=0$



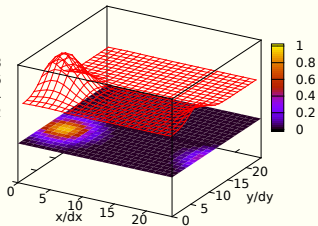
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mpdata<3> $t/dt=0$

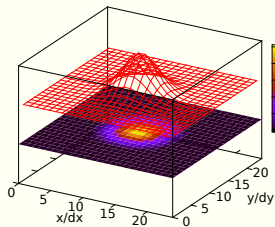


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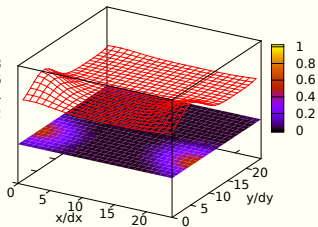


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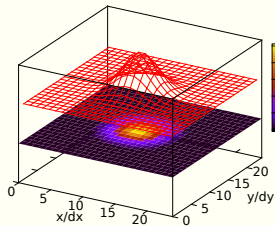
donorcell $t/dt=0$



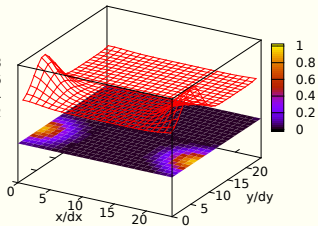
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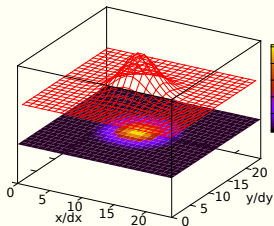


mpdata<3> $t/dt=24$

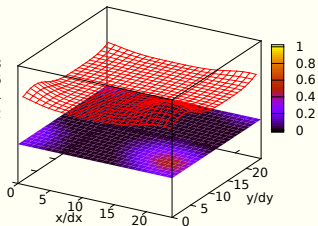


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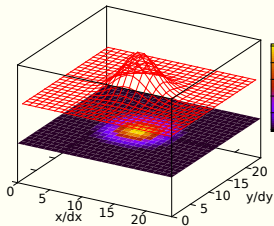
donorcell $t/dt=0$



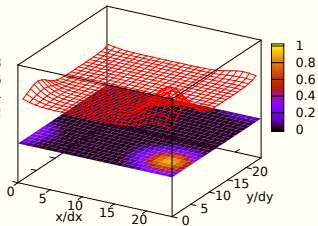
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mpdata<3> $t/dt=0$

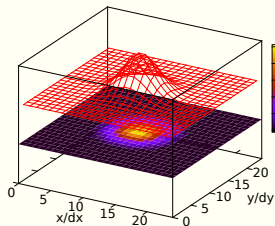


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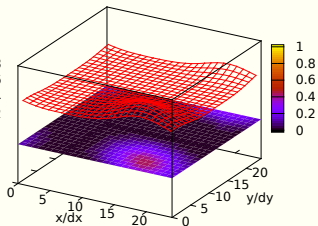


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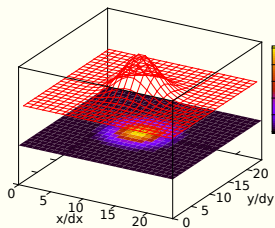
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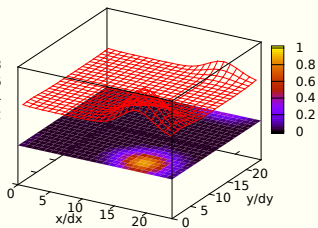
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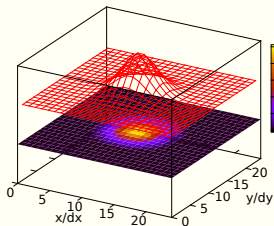


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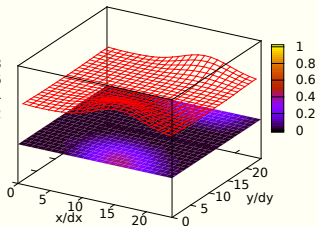


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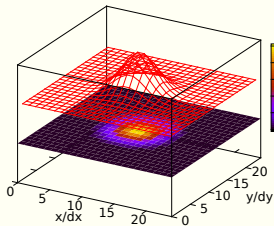
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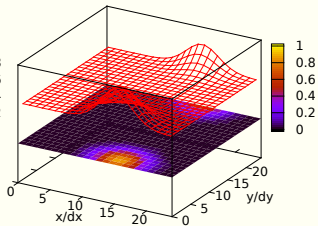
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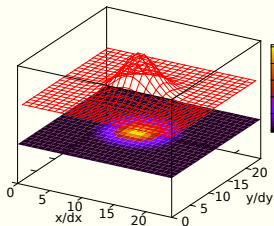


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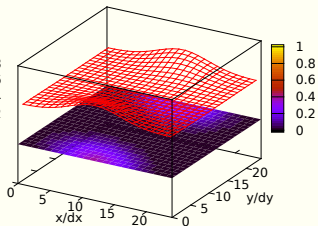


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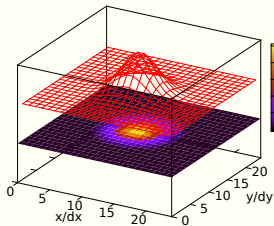
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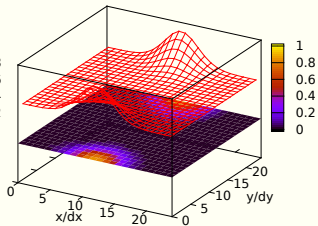
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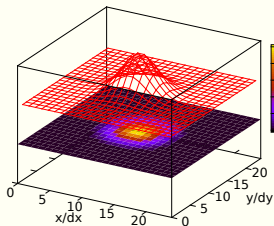


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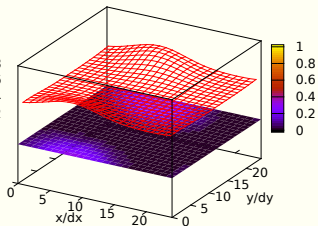


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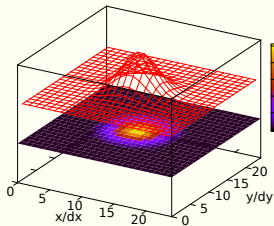
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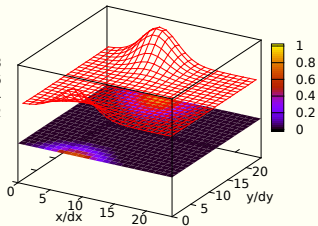
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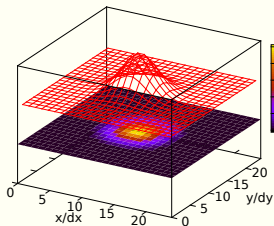


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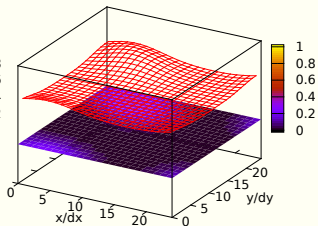


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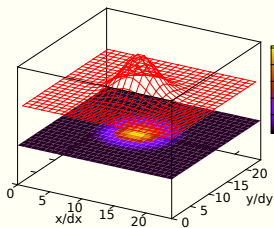
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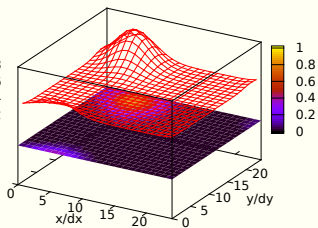
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mpdata<3> $t/dt=0$

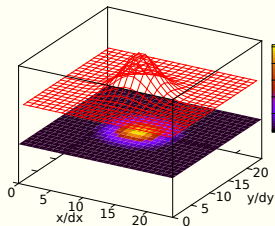


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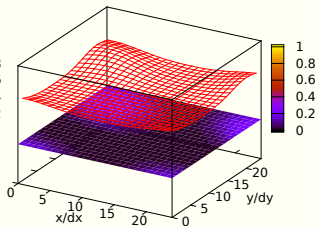


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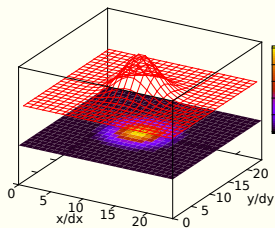
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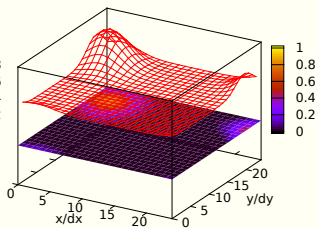
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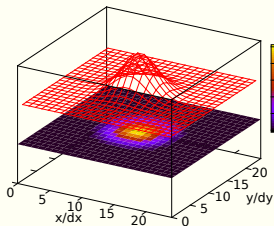


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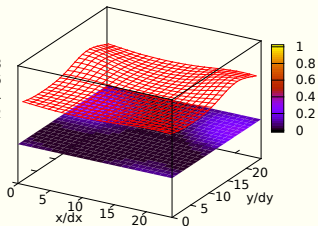


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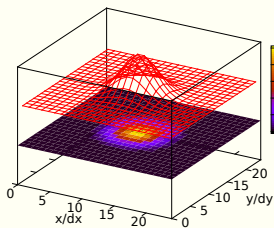
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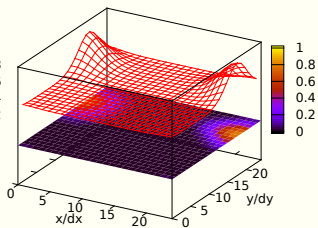
donorcell $t/dt=72$



mpdata<3> $t/dt=0$

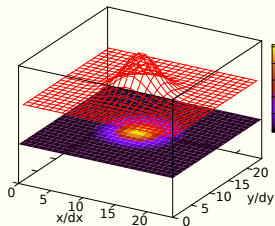


mpdata<3> $t/dt=72$

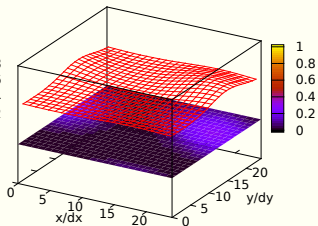


przykład 2D (Arabas et al. 2014, Sci. Prog.)

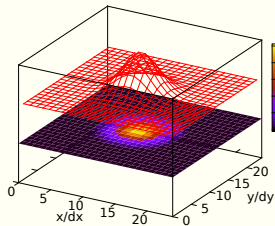
donorcell $t/dt=0$



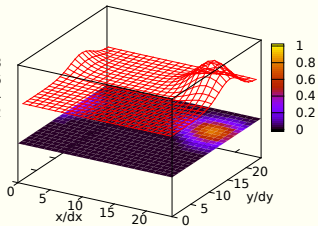
donorcell $t/dt=78$



mpdata<3> $t/dt=0$

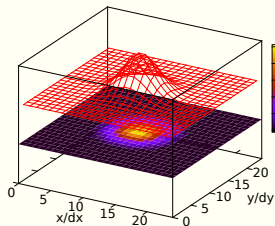


mpdata<3> $t/dt=78$

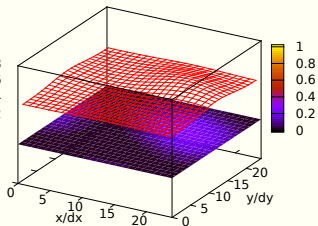


przykład 2D (Arabas et al. 2014, Sci. Prog.)

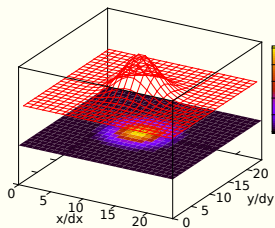
donorcell $t/dt=0$



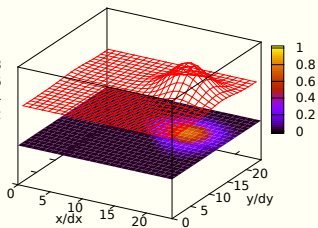
donorcell $t/dt=84$



mpdata<3> $t/dt=0$

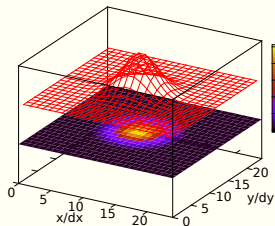


mpdata<3> $t/dt=84$

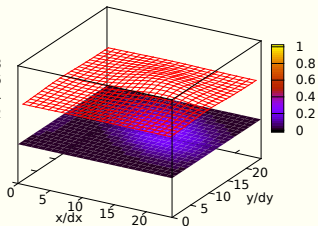


przykład 2D (Arabas et al. 2014, Sci. Prog.)

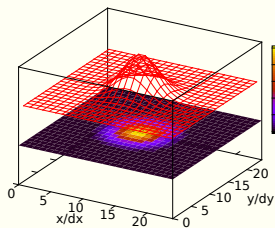
donorcell $t/dt=0$



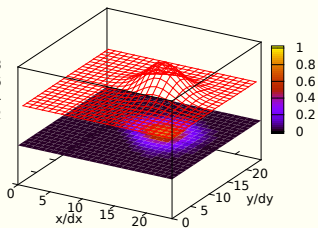
donorcell $t/dt=90$



mpdata<3> $t/dt=0$

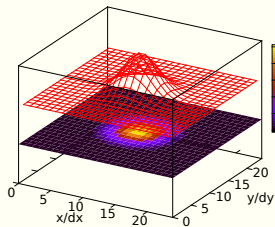


mpdata<3> $t/dt=90$

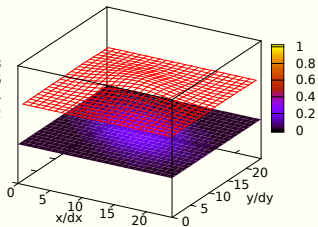


przykład 2D (Arabas et al. 2014, Sci. Prog.)

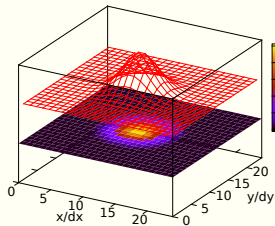
donorcell $t/dt=0$



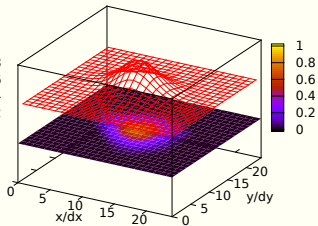
donorcell $t/dt=96$



mpdata<3> $t/dt=0$



mpdata<3> $t/dt=96$



libmpdata++

Jaruga et al. 2015

Geosci. Model Dev., 8, 1005–1032, 2015

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doi:10.5194/gmd-8-1005-2015

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Model Development

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libmpdata++ 1.0: a library of parallel MPDATA solvers for systems of generalised transport equations

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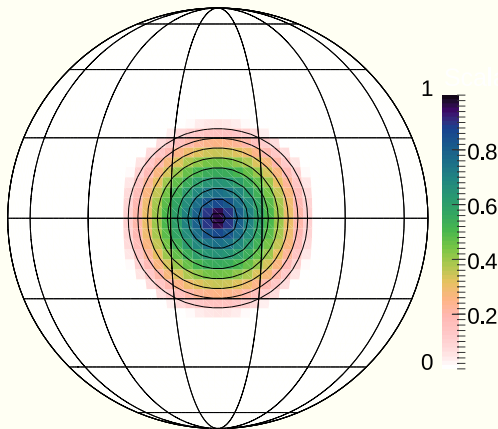
¹Institute of Geophysics, Faculty of Physics, University of Warsaw, Warsaw, Poland

²National Center for Atmospheric Research, Boulder, CO, USA

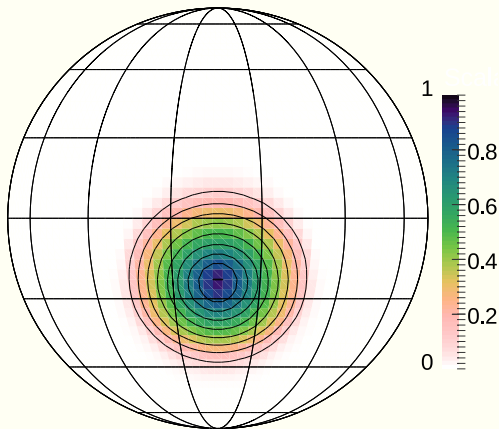
³European Centre for Medium-Range Weather Forecasts, Reading, UK

$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$

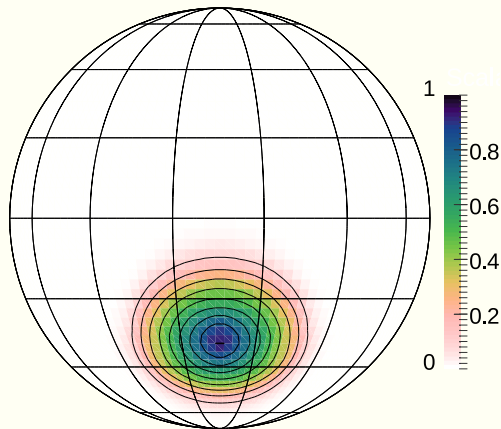
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



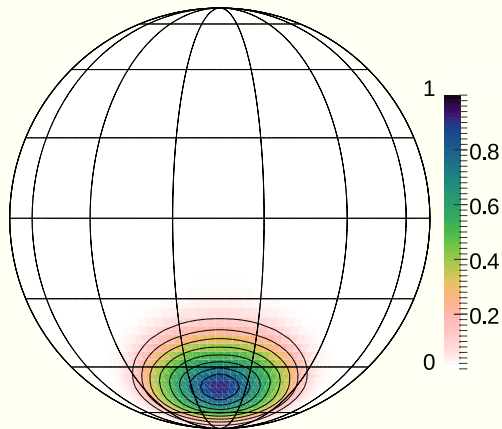
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



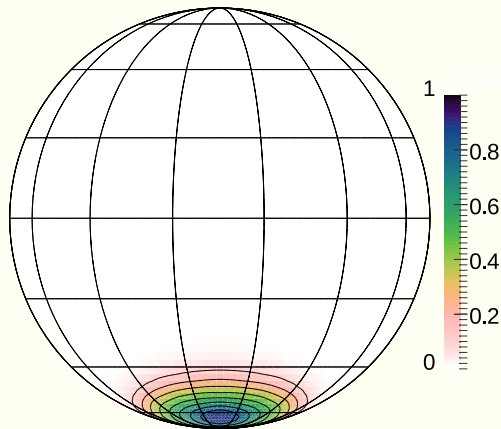
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



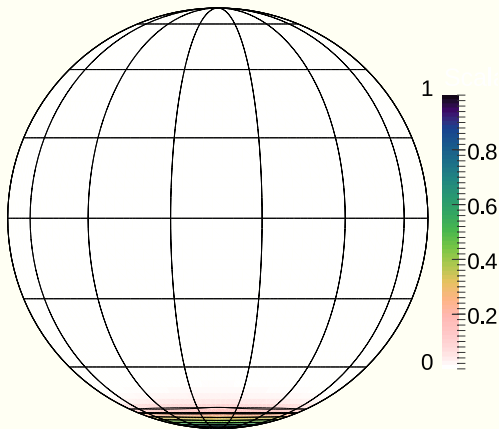
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



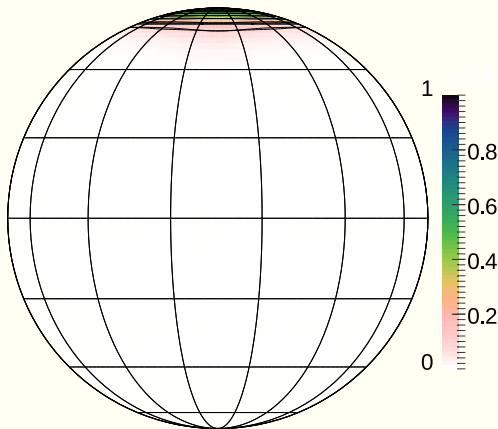
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



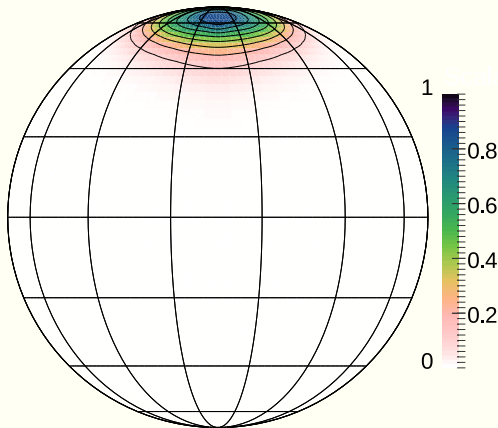
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



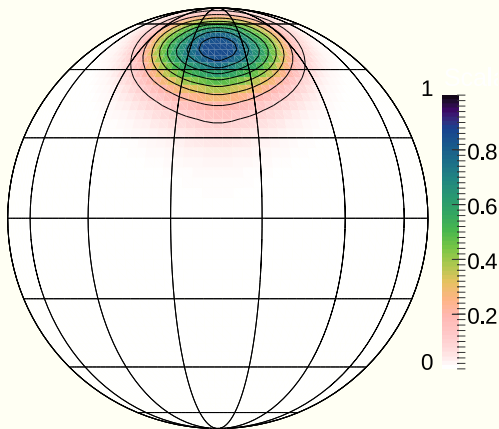
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



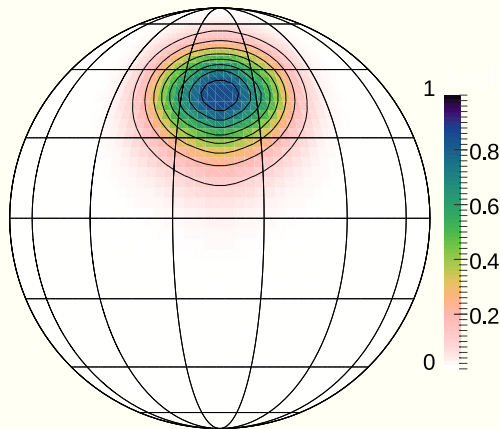
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



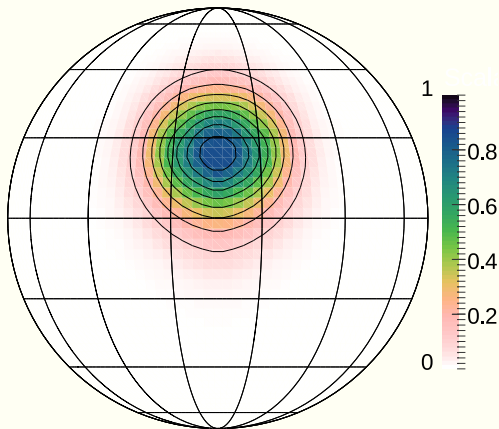
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



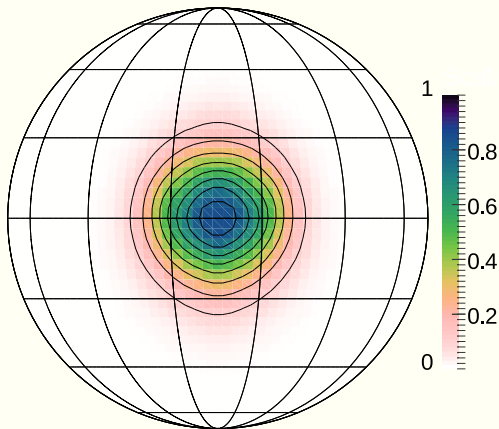
$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



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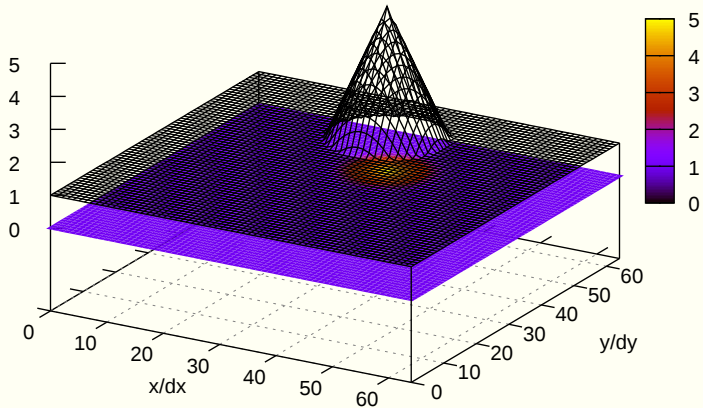


$$\partial_t(G\psi) + \nabla \cdot (G\vec{u}\psi) = GR$$



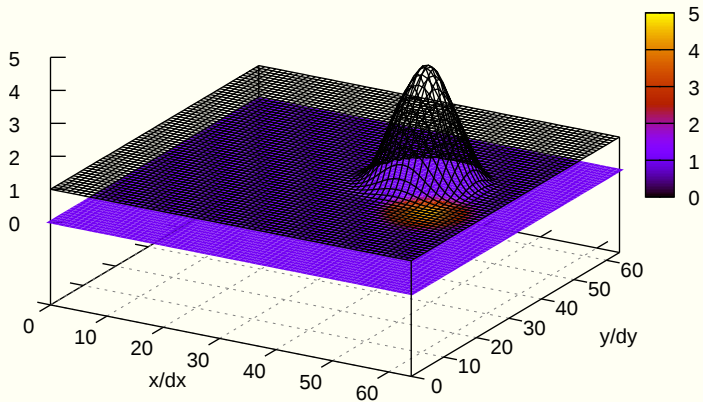
libmpdata++: rotating cone test

($t/dt=0$)



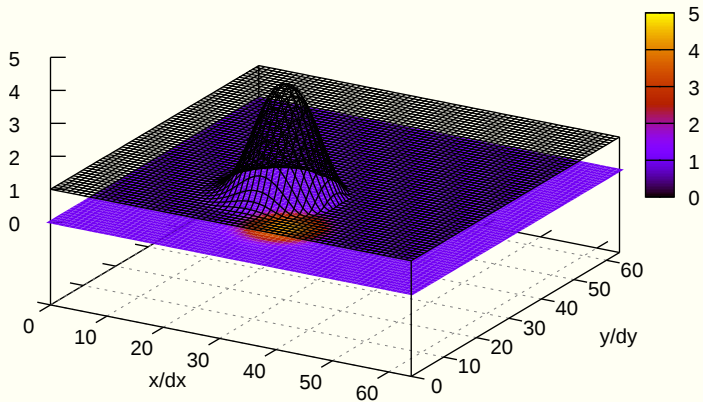
libmpdata++: rotating cone test

($t/dt=157$)



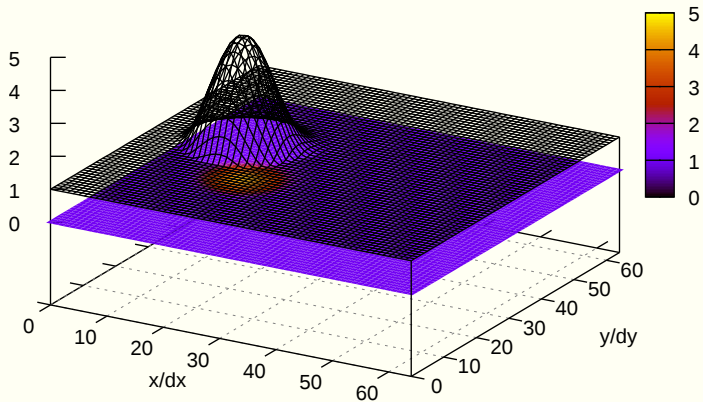
libmpdata++: rotating cone test

(t/dt=314)



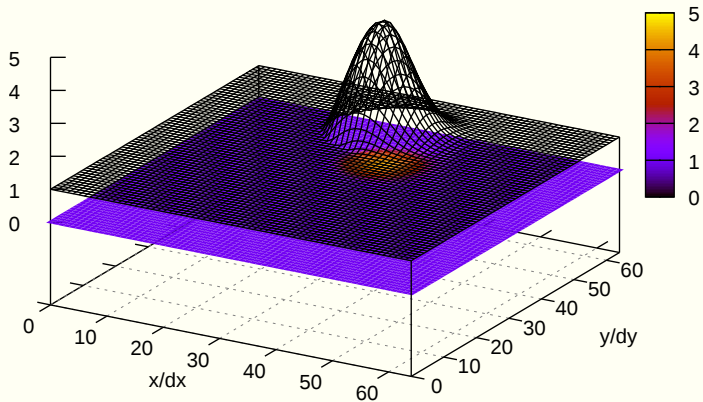
libmpdata++: rotating cone test

(t/dt=471)



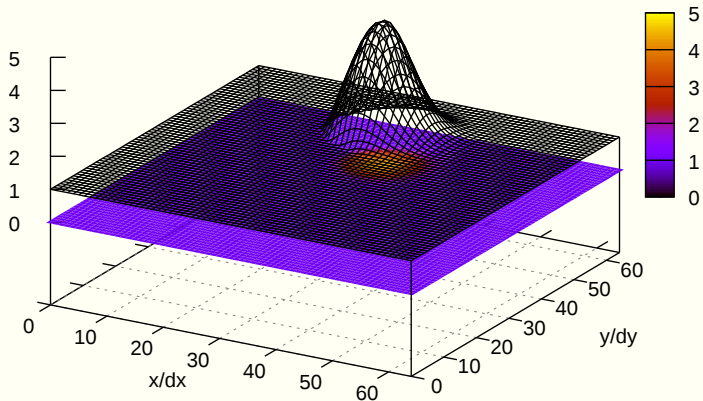
libmpdata++: rotating cone test

($t/dt=628$)



libmpdata++: rotating cone test

(t/dt=628)



64 LOC using libmpdata++

```

1 #include <libmpdata++/solvers/mpdata.hpp>
2 #include <libmpdata++/concurr/serial.hpp>
3 #include <libmpdata++/output/gnuplot.hpp>
4
5 int main()
6 {
7     namespace lmpdt = libmpdataxx;
8     const int nx=64, ny=64, nt = 628;
9
10    // compile-time parameters
11    struct ct_params_t : lmpdt::ct_params_default_t
12    {
13        using real_t = double;
14        enum { n_dims = 2 };
15        enum { n_eqns = 1 };
16    };
17
18    // solver choice
19    using run_t = lmpdt::output::gnuplot< lmpdt::solvers::mpdata< ct_params_t >>;
20
21    // runtime parameters
22    typename run_t::rt_params_t p;
23    p.grid_size = {nx+1, ny+1};
24    p.outfreq = nt/4;
25    p.gnuplot_output = "out_%s_%d.svg";
26    p.gnuplot_with = "lines";
27    p.gnuplot_cbrange = p.gnuplot_zrange = "[0:5]";
28
29    // sharedmem concurency and boundary condition choice
30    lmpdt::concurr::serial<
31        run_t,
32        lmpdt::bcond::open, lmpdt::bcond::open, // x-left, x-right
33        lmpdt::bcond::open, lmpdt::bcond::open // y-left, y-right
34    > run(p);

```



```

35
36 // initial condition
37 {
38     using namespace blitz::tensor;
39     auto psi = run.advectee();
40
41     const double
42         dt = .1, dx = 1, dy = 1, omega = .1,
43         h = 4., h0 = 1, r = .15 * nx * dx,
44         x0 = .5 * nx * dx, y0 = .75 * ny * dy,
45         xc = .5 * nx * dx, yc = .50 * ny * dy;
46
47     // cone shape cut at h0
48     psi = blitz::pow(i * dx - x0, 2) +
49           blitz::pow(j * dy - y0, 2);
50
51     psi = h0 + where(
52         psi - pow(r, 2) <= 0,           // if
53         h - blitz::sqrt(psi / pow(r/h,2)), // then
54         0.                             // else
55     );
56
57     // constant-angular-velocity rotational field
58     run.advector(0) = omega * (j * dy - yc) * dt/dx;
59     run.advector(1) = -omega * (i * dx - xc) * dt/dy;
60 }
61
62 // time stepping
63 run.advance(nt);
64 }

```

```

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37 {
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39     auto psi = run.advectee();
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61
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63 run.advance(nt);
64 }

```

CMakeLists.txt

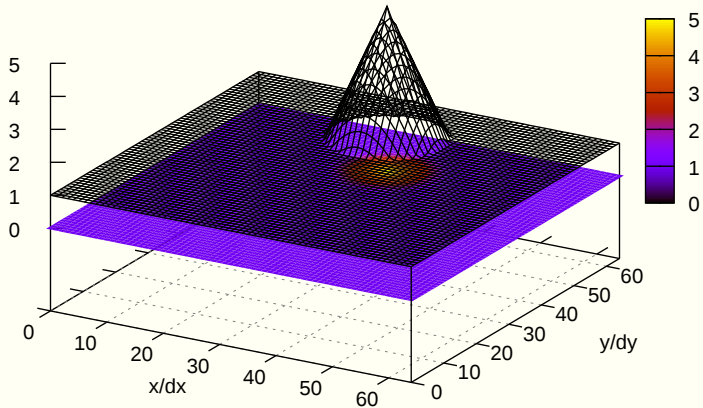
```

1 cmake_minimum_required(VERSION 3.0)
2 project(hello_world CXX)
3 find_package(libmpdata++)
4 set(CMAKE_CXX_FLAGS ${libmpdataxx_CXX_FLAGS_RELEASE})
5 add_executable(hello_world hello_world.cpp)
6 target_link_libraries(hello_world ${libmpdataxx_LIBRARIES})

```

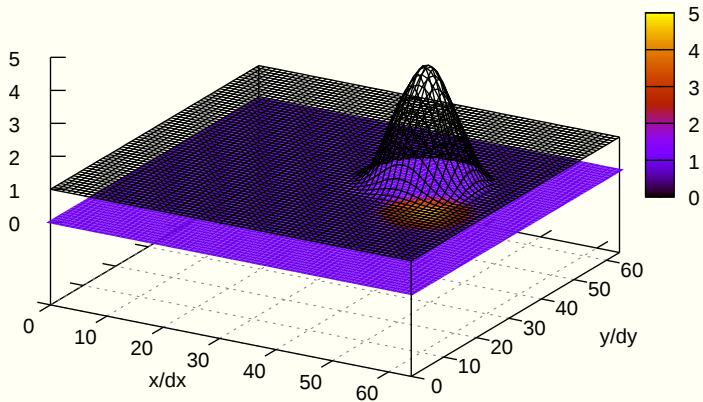
libmpdata++: rotating cone test

($t/dt=0$)



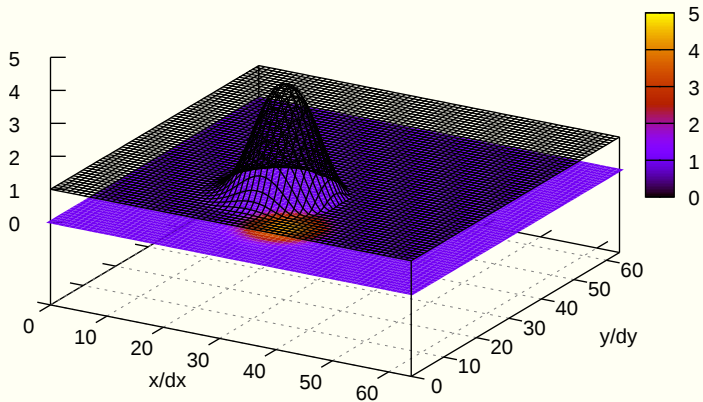
libmpdata++: rotating cone test

($t/dt=157$)



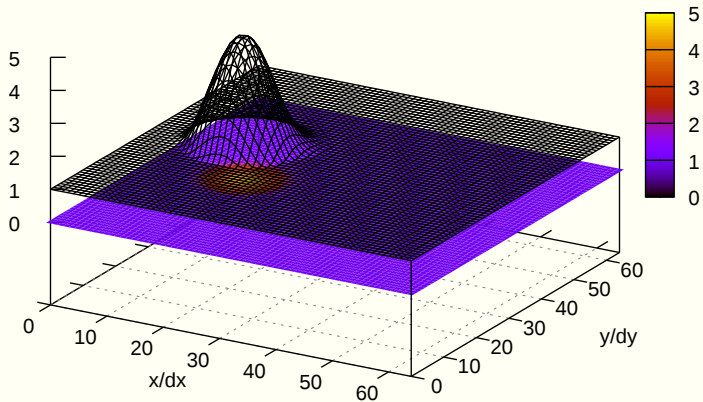
libmpdata++: rotating cone test

(t/dt=314)



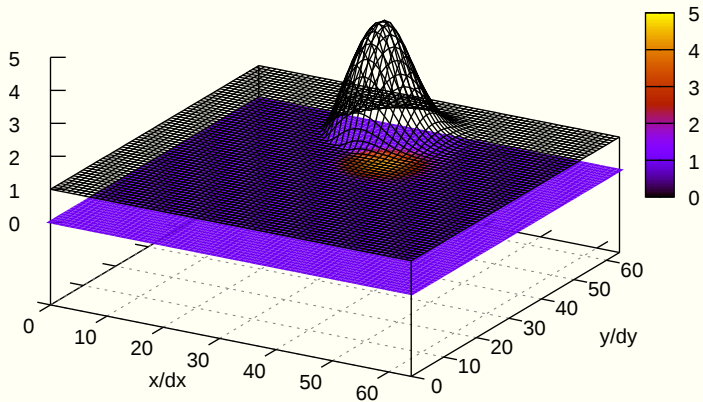
libmpdata++: rotating cone test

(t/dt=471)



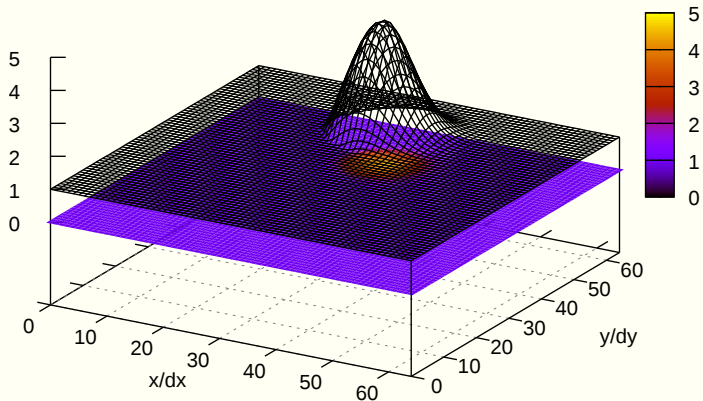
libmpdata++: rotating cone test

($t/dt=628$)



libmpdata++: rotating cone test

(t/dt=628)



64 LOC using libmpdata++

with multi-threading \rightsquigarrow also 64 LOC!

```
2c2
< #include <libmpdata++/concurr/serial.hpp>
---
> #include <libmpdata++/concurr/threads.hpp>
30c30
<   lmpdt::concurr::serial<
---
>   lmpdt::concurr::threads<
```

```
$ top
```

```
...
  PID USER      PR  NI  S   %CPU %MEM  nTH      TIME+  COMMAND  %MEM
21031 slayoo    20   0   R  73.7  0.1    4    0:01.68  hello_worl  90%
...
```

MPI + threads \rightsquigarrow also 64 LOC!!! (recompilation only)

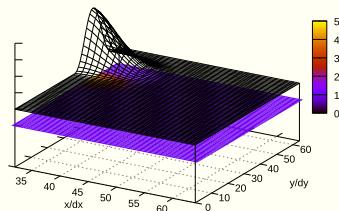
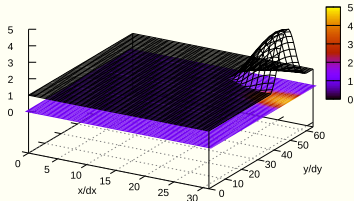
```
$ cmake . -DCMAKE_CXX_COMPILER=mpic++  
$ make  
$ OMP_NUM_THREADS=2 mpirun -np 2 ./hello_world
```

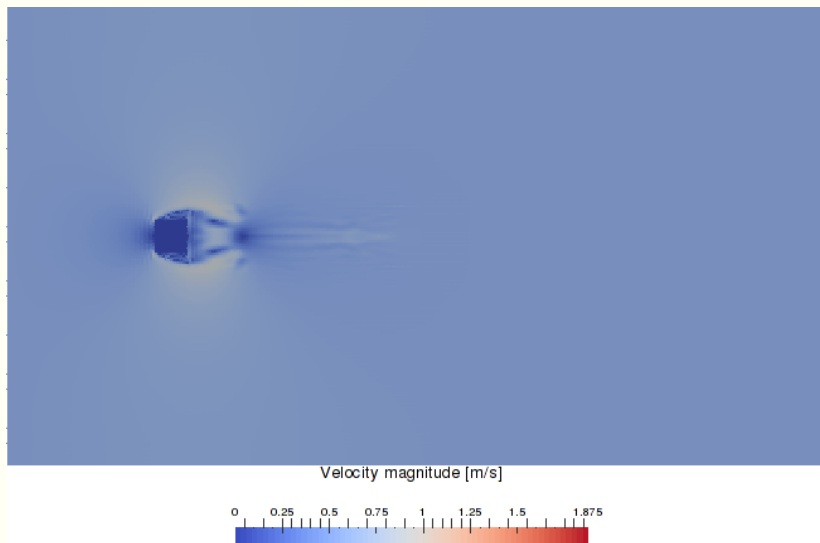
```
$ top
```

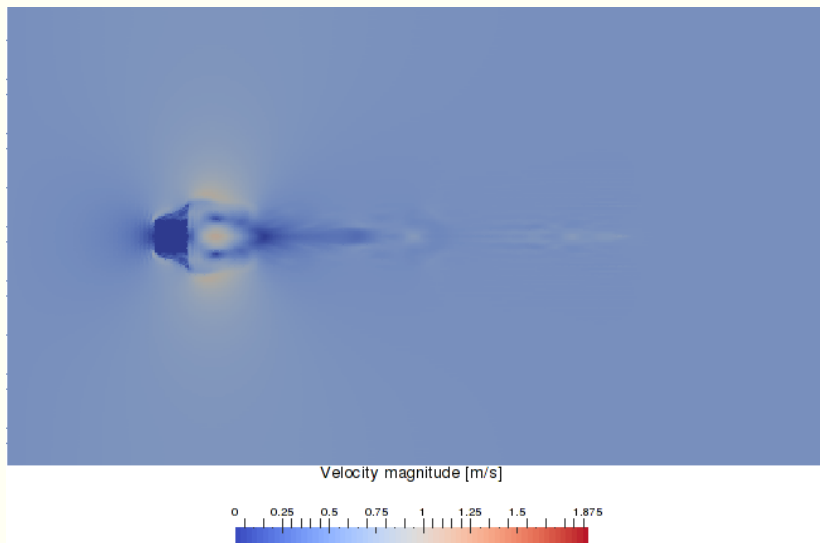
```
...
```

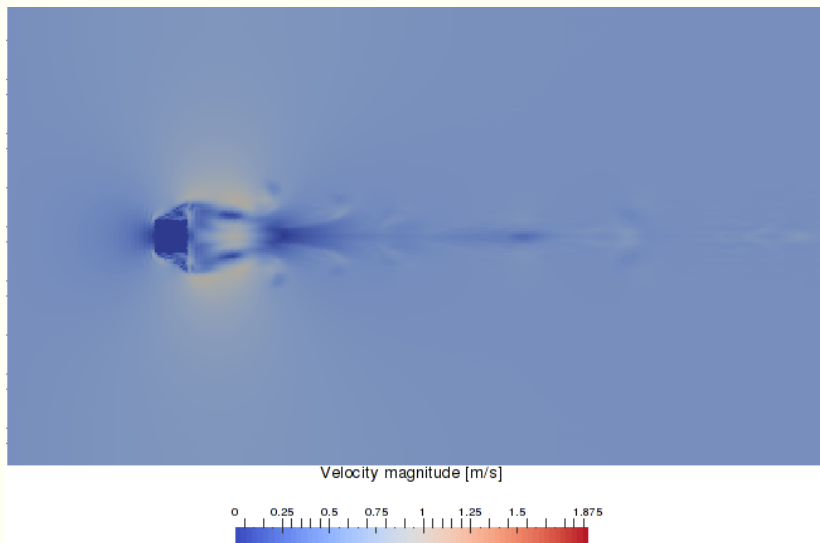
PID	USER	PR	NI	S	%CPU	%MEM	nTH	TIME+	COMMAND	
19640	slayoo	20	0	R	65.5	0.3	2	0:00.92	hello_worl	98%
19641	slayoo	20	0	R	64.0	0.3	2	0:00.91	hello_worl	99%

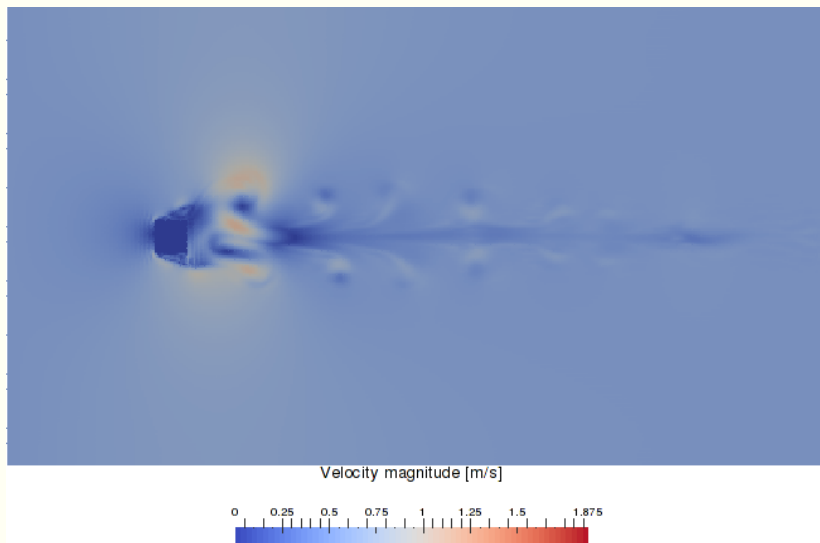
```
...
```

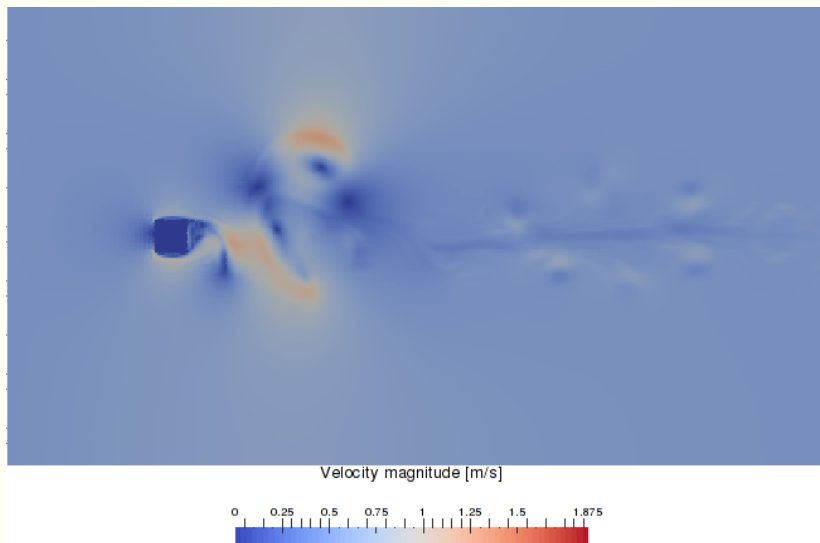


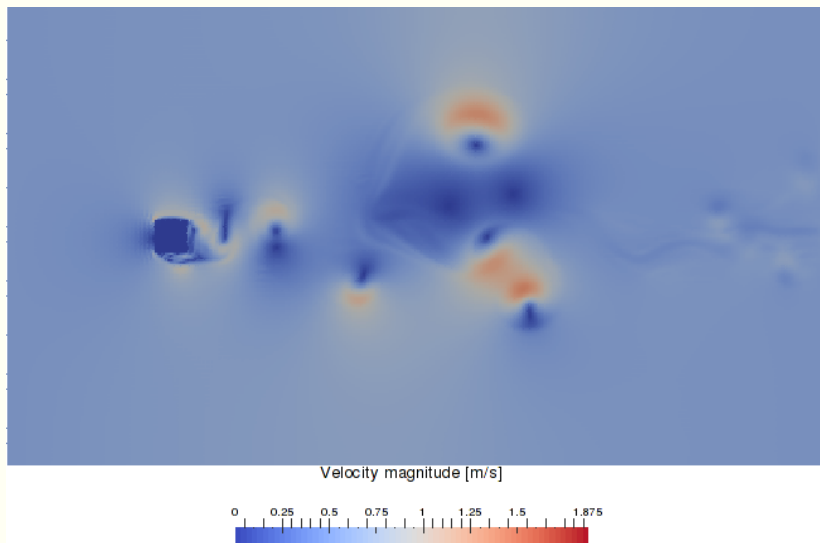


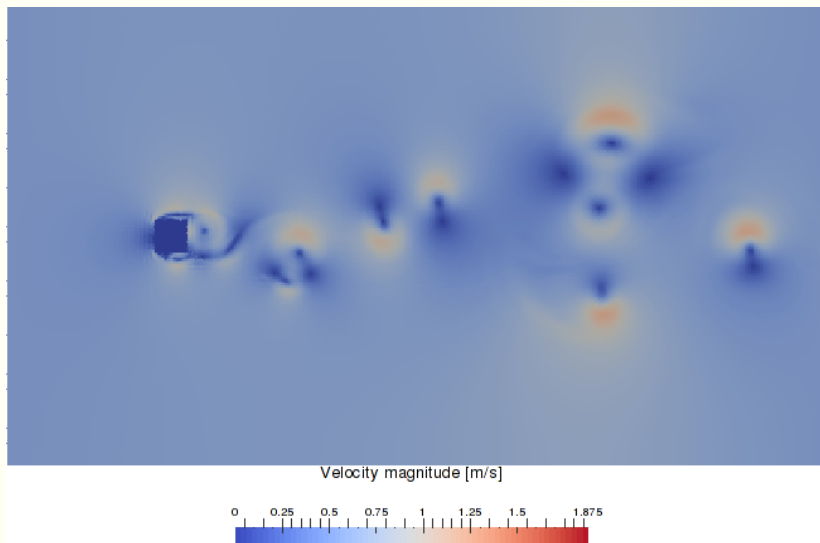


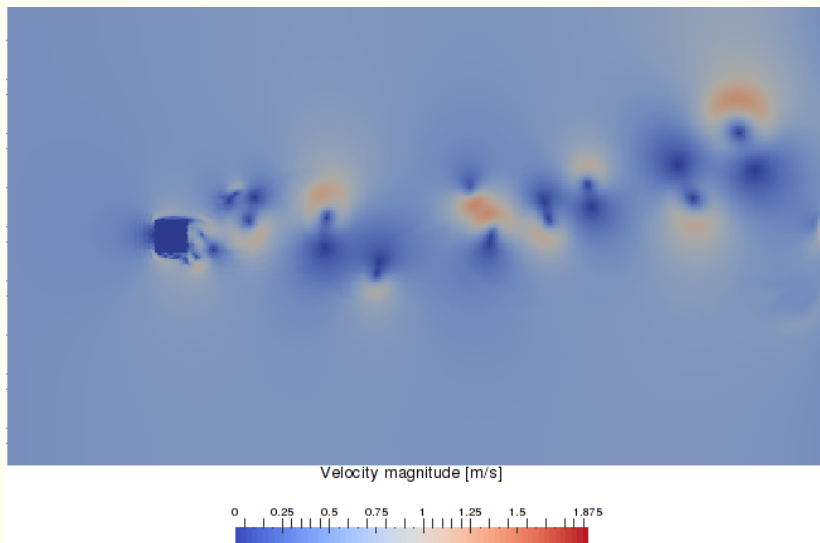


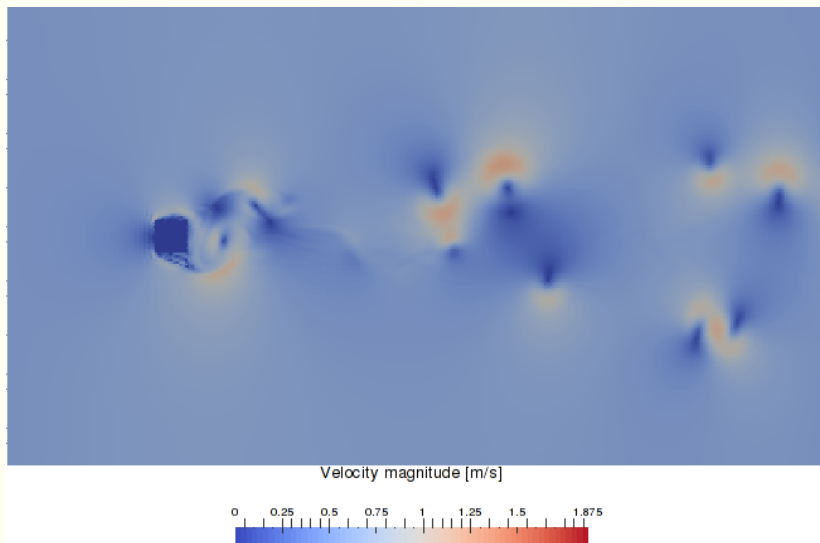


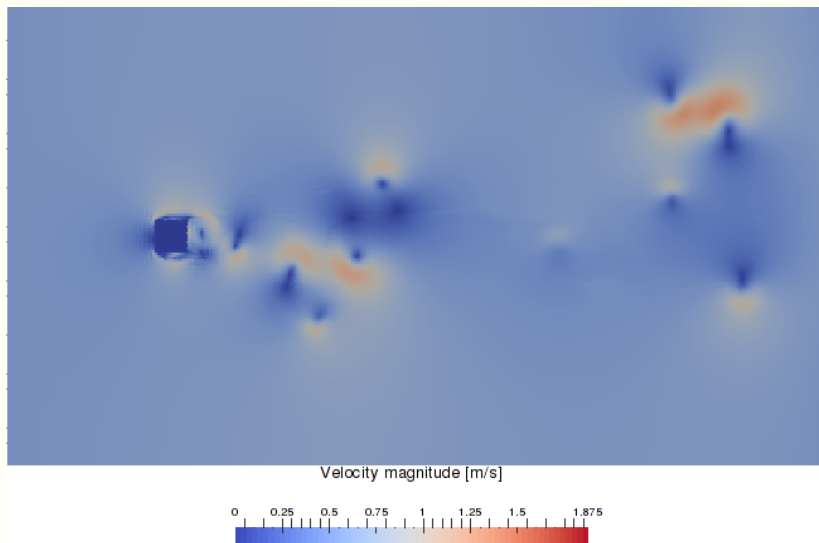


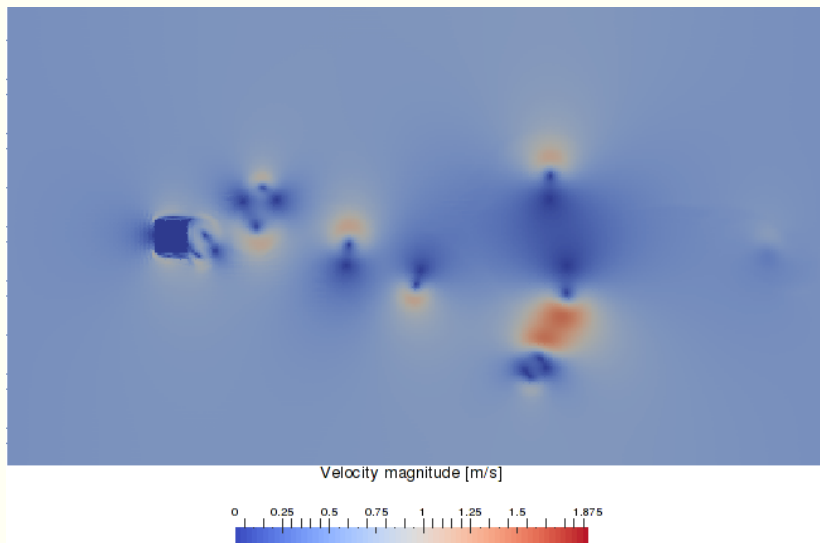


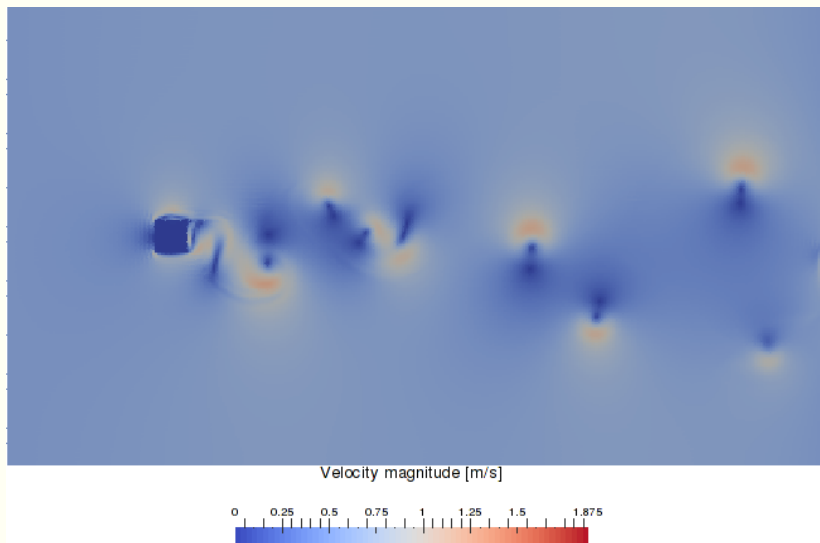


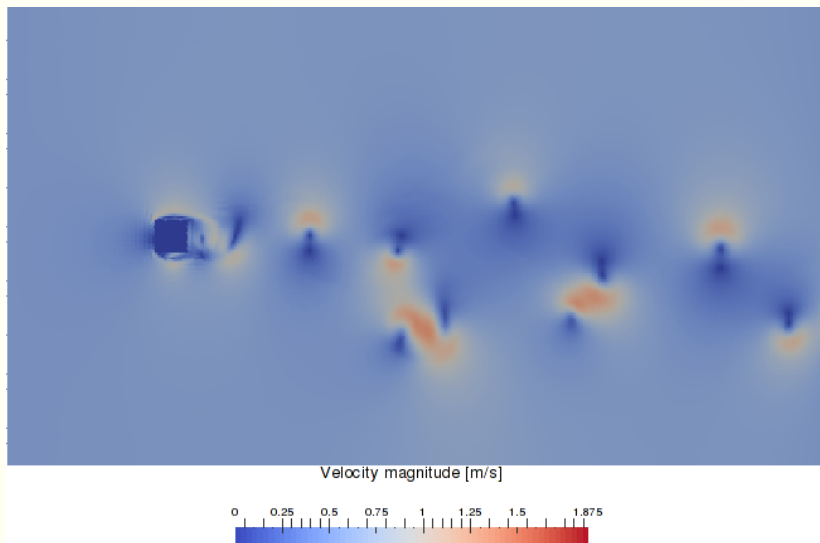


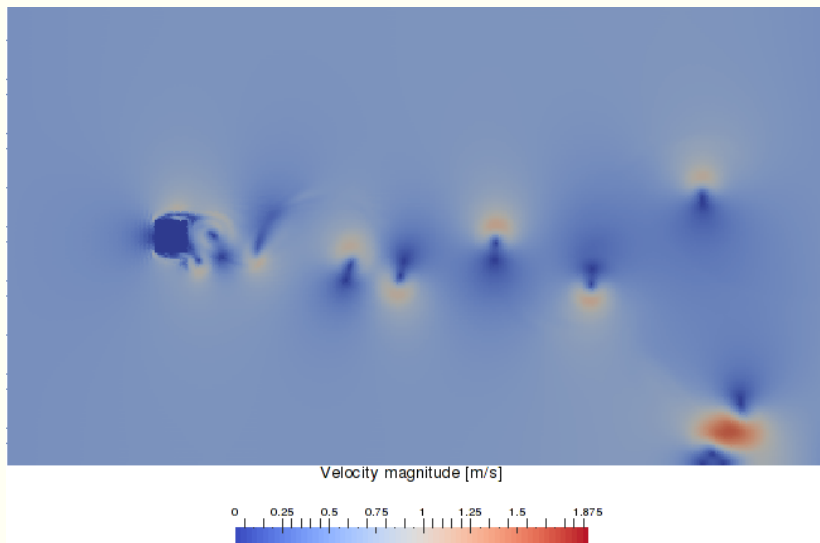


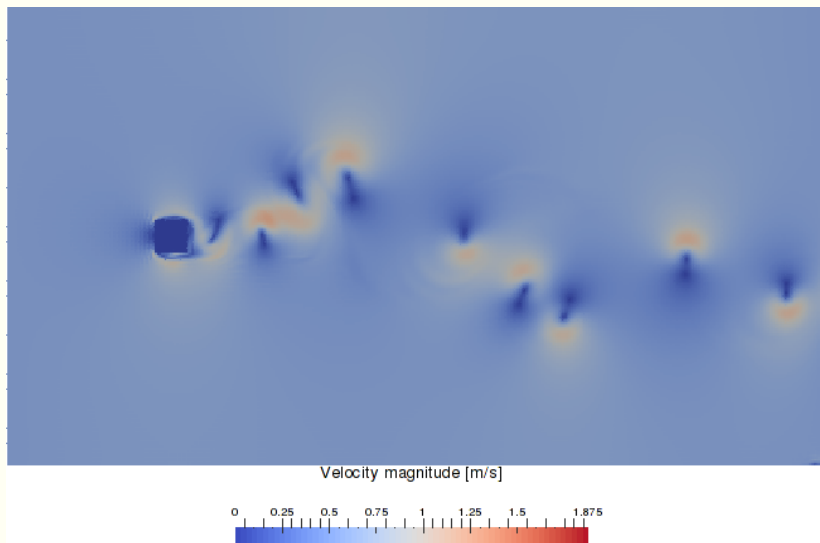


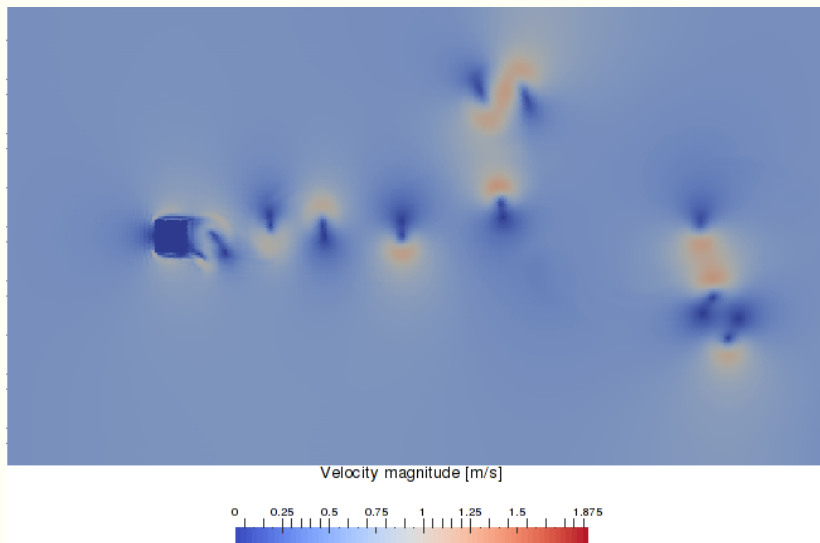


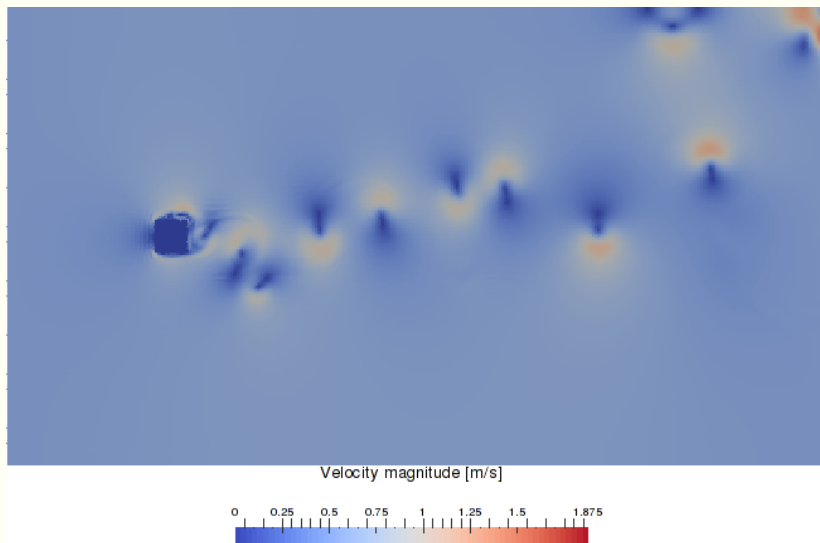


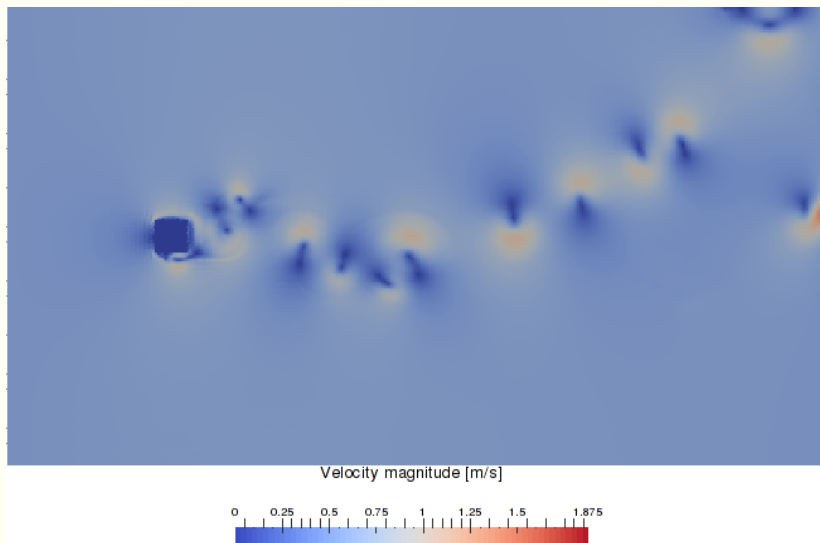


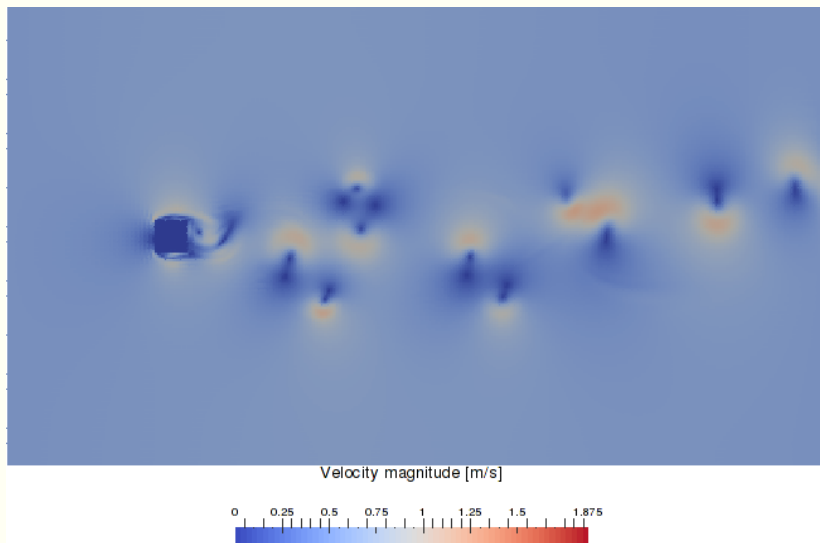


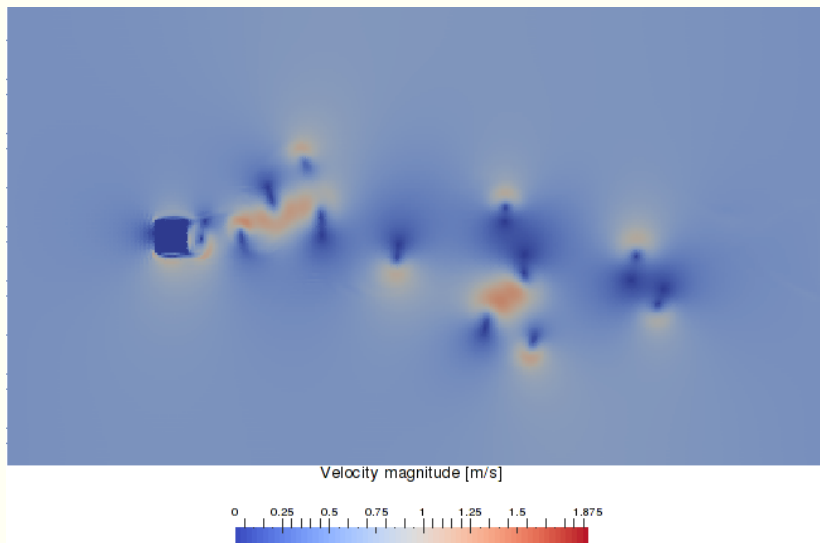


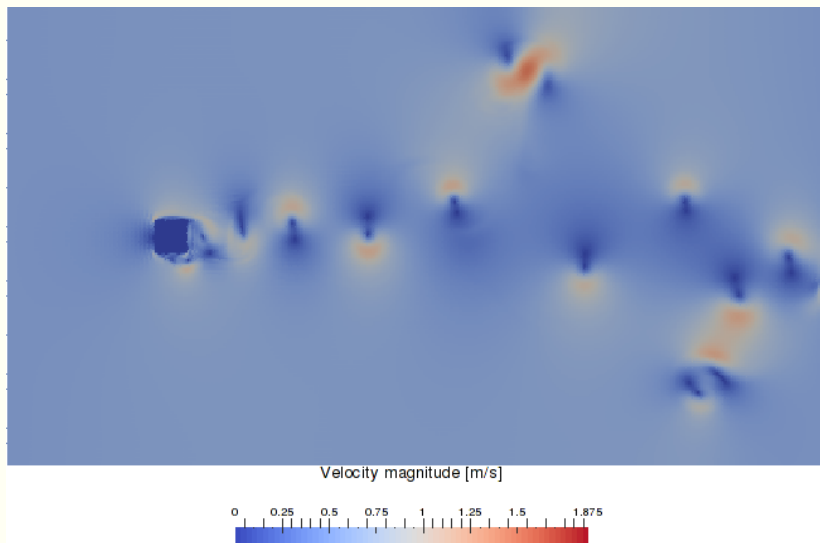


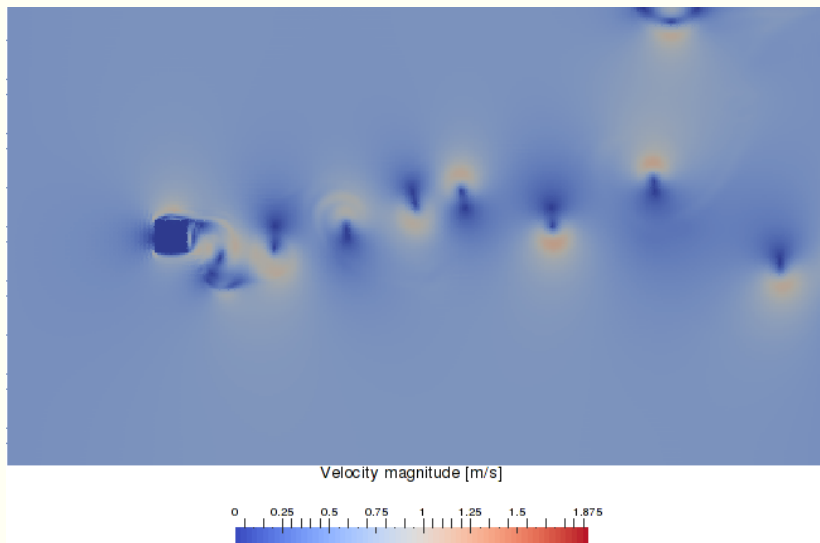


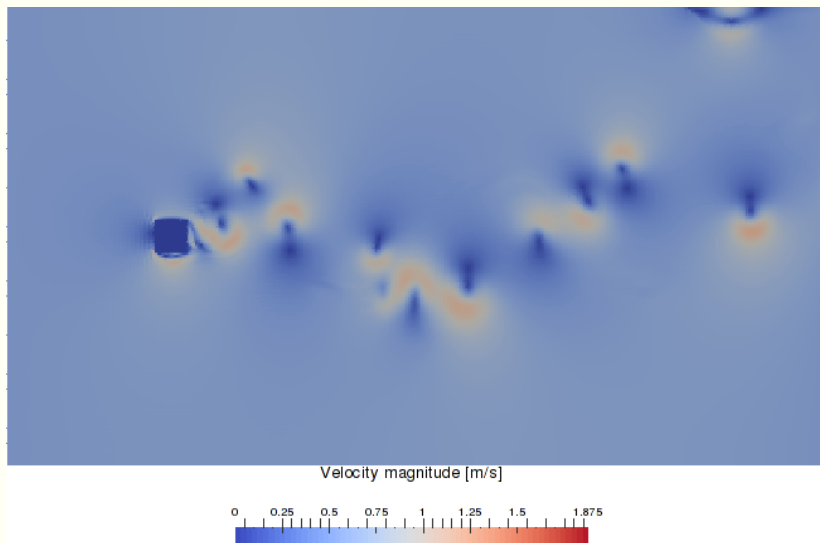


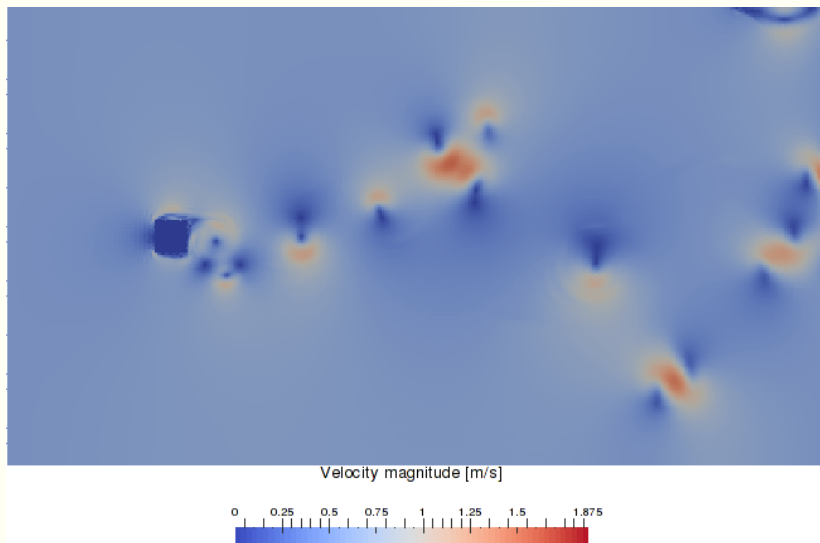


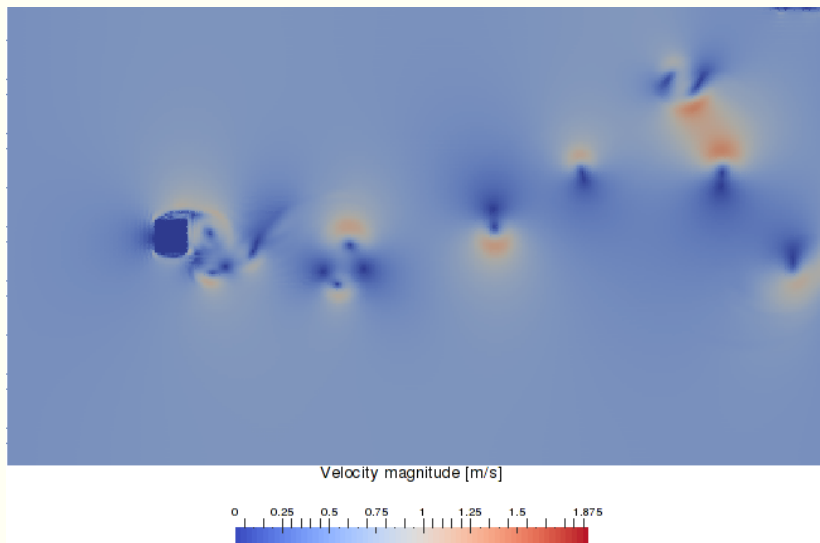


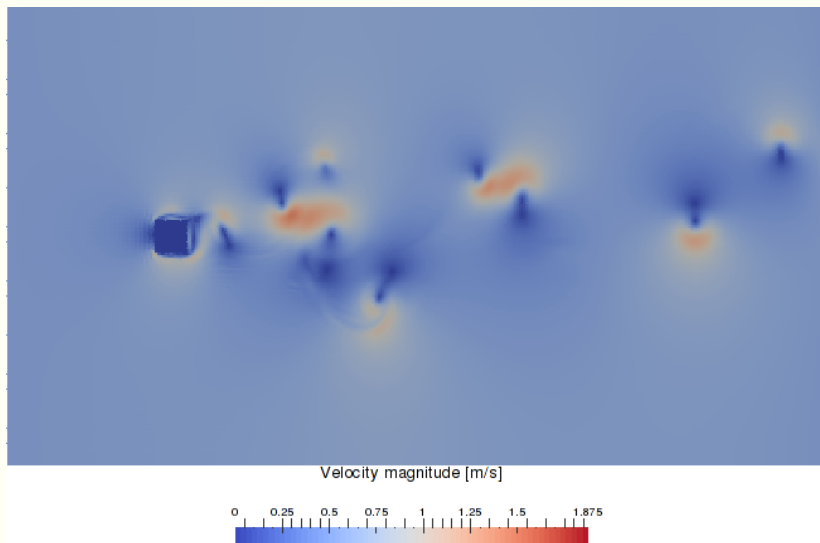


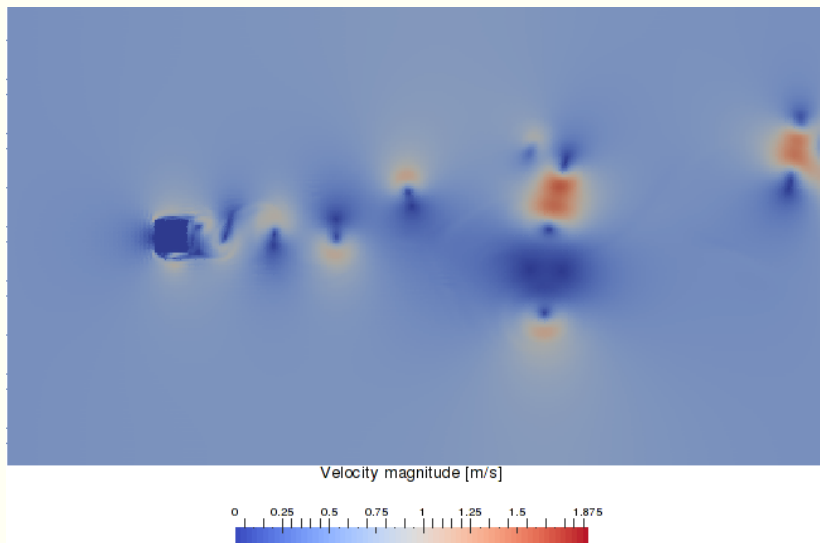


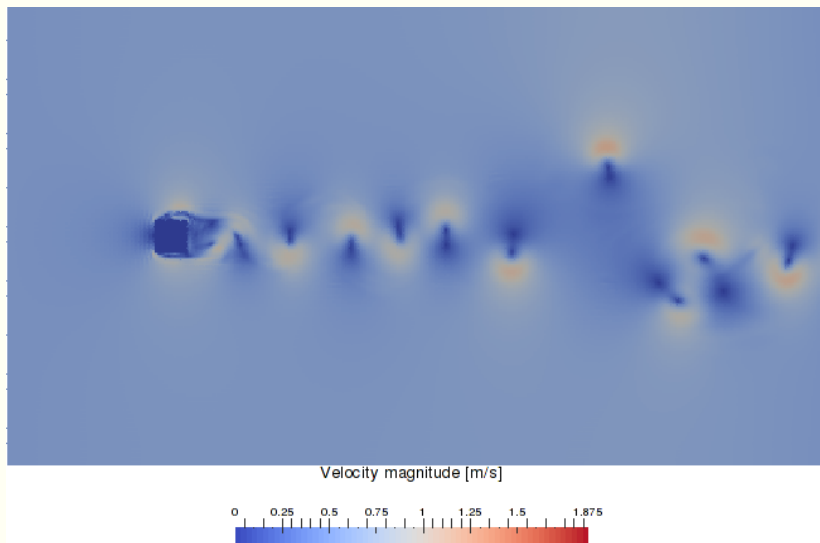


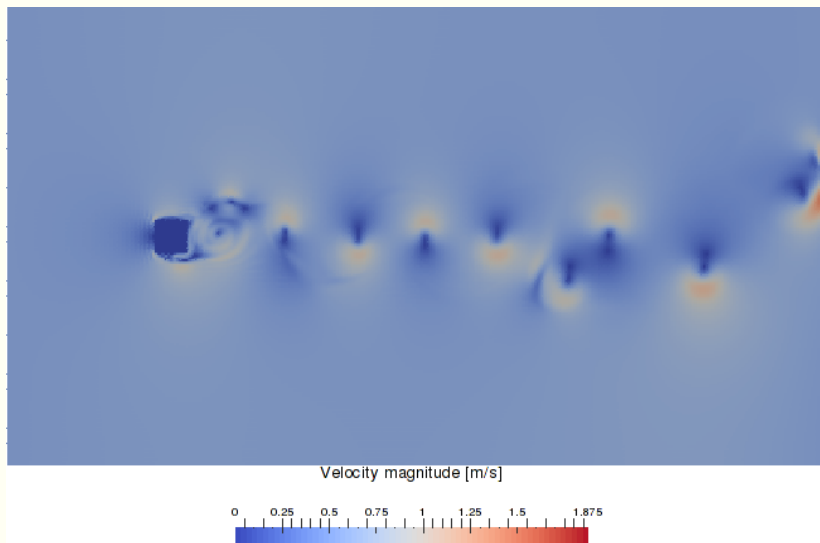


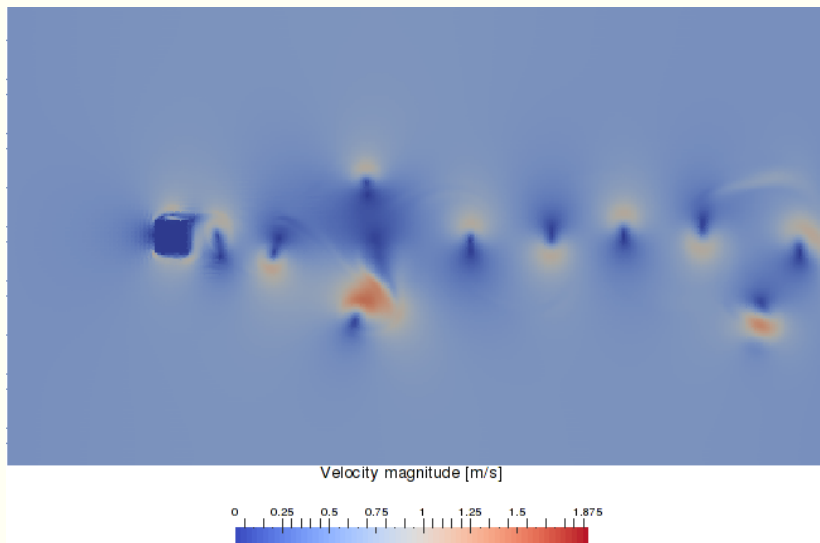


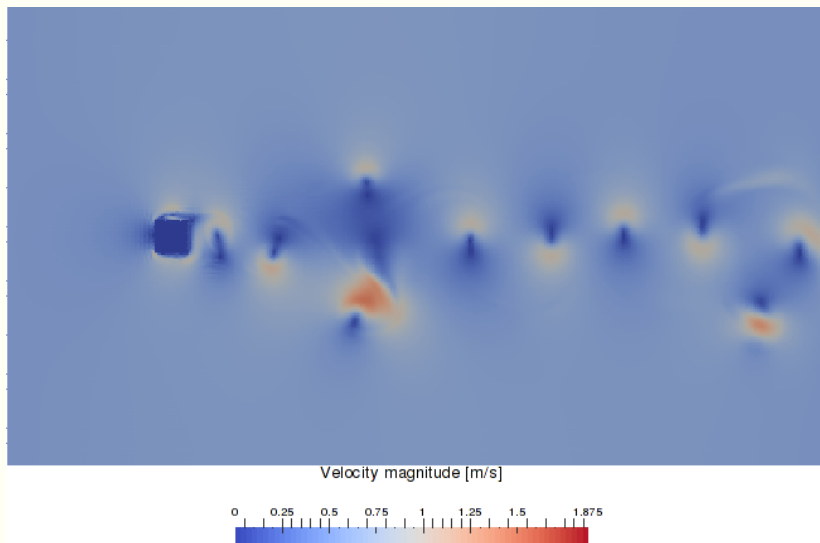












<https://www.youtube.com/watch?v=BEidkhpw-MA>

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key features:

- reusable – API documented in the paper; out-of-tree setups
- comprehensive set of MPDATA opts (incl. FCT, infinite-gauge, ...)
- 1D, 2D & 3D integration; optional coordinate transformation
- four types of solvers:
 - `mpdata::mpdata` (Euler equations)
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 - `mpdata::mpdata3` (Euler equations)
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- parallelisation: threads + MPI
- separation of concerns (numerics / boundary cond. / io / concurrency)
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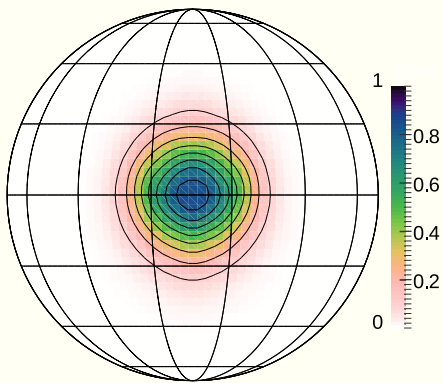
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re last step: Smolarkiewicz and Clark (1986, JCP), Sousa (2009, IJNMF),
Smolarkiewicz and Szmelter (2005, JCP), Cristiani (2015, JCSMD)

same trick!

MPDATA in a nutshell (Smolarkiewicz 1983, 1984, ...)

$$\text{transport PDE: } \frac{\partial \psi}{\partial t} + \frac{\partial}{\partial x} (v\psi) = 0$$

$$\psi_i^{n+1} = \psi_i^n - [F(\psi_i^n, \psi_{i+1}^n, C_{i+1/2}) - F(\psi_{i-1}^n, \psi_i^n, C_{i-1/2})]$$

$$F(\psi_L, \psi_R, C) = \max(C, 0) \cdot \psi_L + \min(C, 0) \cdot \psi_R$$

$$C = v\Delta t / \Delta x$$

upwind

$$\text{modified eq.: } \frac{\partial \psi}{\partial t} + \frac{\partial}{\partial x} (v\psi) + \underbrace{K \frac{\partial^2 \psi}{\partial x^2}}_{\text{numerical diffusion}} + \dots = 0 \quad \leftarrow \text{MEA}$$

$$\frac{\partial \psi}{\partial t} + \frac{\partial}{\partial x} (v\psi) + \frac{\partial}{\partial x} \left[\underbrace{\left(-\frac{K}{\psi} \frac{\partial \psi}{\partial x} \right)}_{\text{antidiffusive flux}} \psi \right] = 0$$

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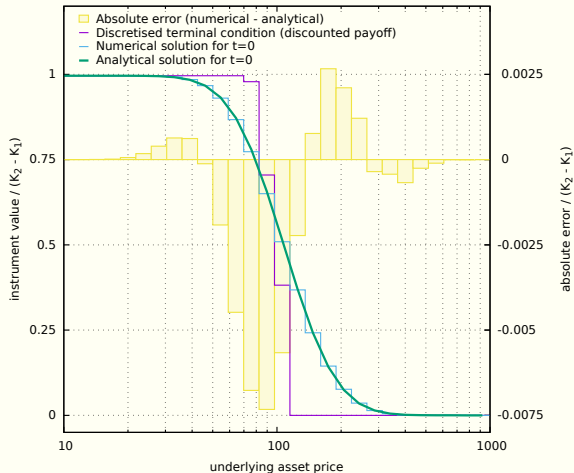
$$\frac{\partial \psi}{\partial t} + \frac{\partial}{\partial x} \left[\left(u - \frac{\nu}{\psi} \frac{\partial \psi}{\partial x} \right) \psi \right] = 0$$

MPDATA meets Black-Scholes: test case

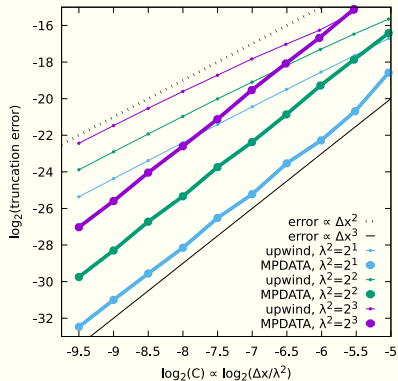
■ payoff function:
corridor

■ truncation error est.
(ψ_a : B-S formula):

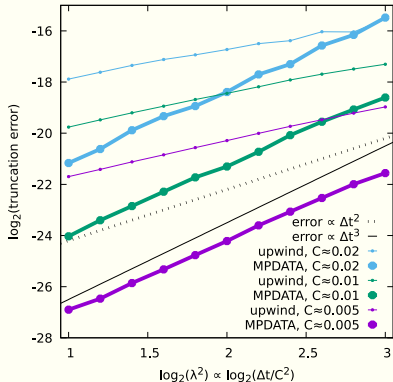
$$E = \sqrt{\sum_{i=1}^{n_x} [\psi_n(x_i) - \psi_a(x_i)]^2 / (n_x \cdot n_t)} \Big|_{t=0}$$



MPDATA meets Black-Scholes: convergence analysis



Truncation error as a function of the Courant number $C = u \frac{\Delta t}{\Delta x}$ which, for fixed λ^2 , is proportional to the gridstep.



Truncation error as a function of the λ^2 parameter which, for fixed C , is proportional to the timestep.

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Dziękuję za uwagę!